

# Unstable periodic orbits and scars in the Dicke model

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Virtual Gathering 2021  
Ergodicity and chaos in  
many-body systems

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# Bohr-Sommerfeld (Regular systems)

Periodic orbits in classical limit



Quantum spectrum ( $\epsilon_i$ )

$$\oint p dq = n(2\pi\hbar) \quad \begin{array}{c} \text{Quantization} \\ \leftarrow \rightleftarrows \rightarrow \end{array} \quad \int_{\epsilon_i}^{\epsilon_{i+1}} T d\epsilon = 2\pi\hbar$$

# Example: The harmonic oscillator $\hat{H} = \hbar\omega \left( \hat{a}^\dagger \hat{a} + \frac{1}{2} \right)$

Classical Hamiltonian

$$H = \frac{\omega}{2} \left( q^2 + p^2 + \frac{1}{2} \right)$$

Eigenenergies  $\epsilon_n$

$$\int_{\epsilon_{n+1}}^{\epsilon_n} T d\epsilon = 2\pi\hbar$$

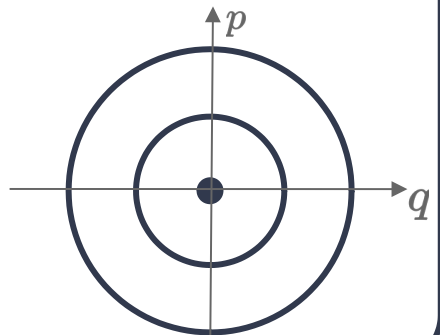
$$\Rightarrow \epsilon_{n+1} - \epsilon_n = \boxed{\hbar\omega}$$

Periodic orbits of energy  $\epsilon$

$$\mathcal{O}_\epsilon(t) = \left( \underbrace{A(\epsilon) \sin(\omega t)}_{q(t)}, \underbrace{A(\epsilon) \cos(\omega t)}_{p(t)} \right)$$

with period

$$T = \frac{2\pi}{\omega}$$



# Regular Systems

Periodic orbits in classical limit

Quantization

Quantum  
eigenenergies

Quantum  
eigenstates?

# Example: The harmonic oscillator $\hat{H} = \omega \left( \hat{a}^\dagger \hat{a} + \frac{1}{2} \right)$

**Husimi function**

$$Q_n(q, p) = |\langle q, p | n \rangle|^2$$

Allows to visualize how state  $|n\rangle$  is distributed in phase space.

$n=0$

$n=1$

$n=5$

$n=10$

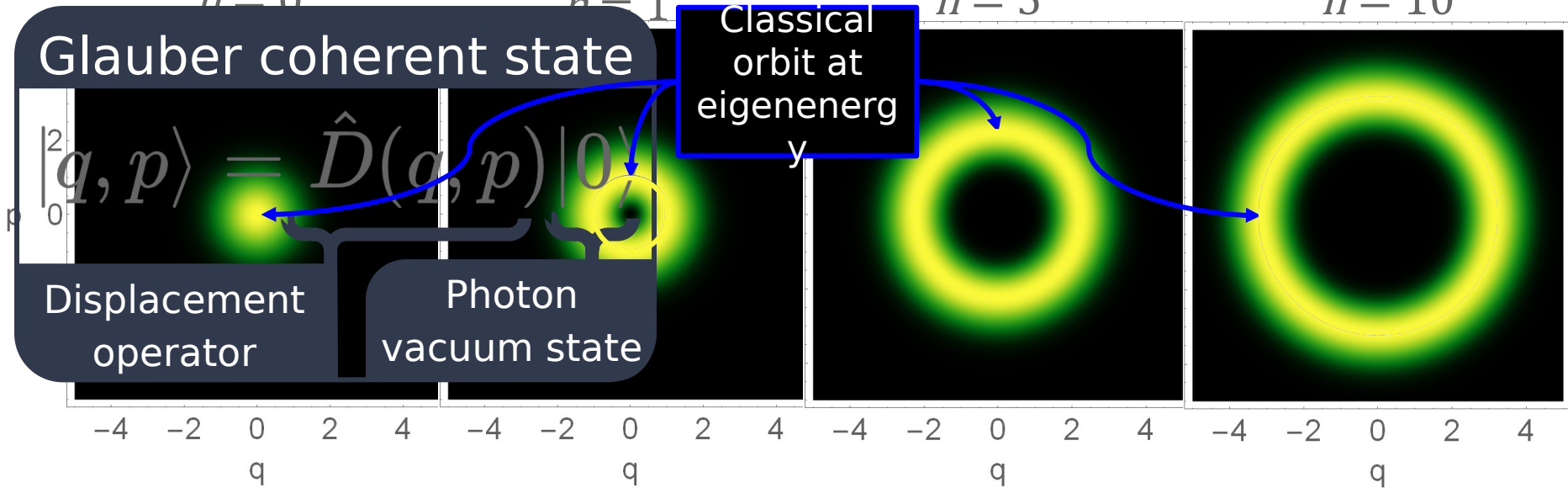
Glauber coherent state

$$|q, p\rangle = \hat{D}(q, p) |0\rangle$$

Displacement operator

Photon vacuum state

Classical orbit at eigenenergy



# Regular Systems

Periodic orbits in classical limit

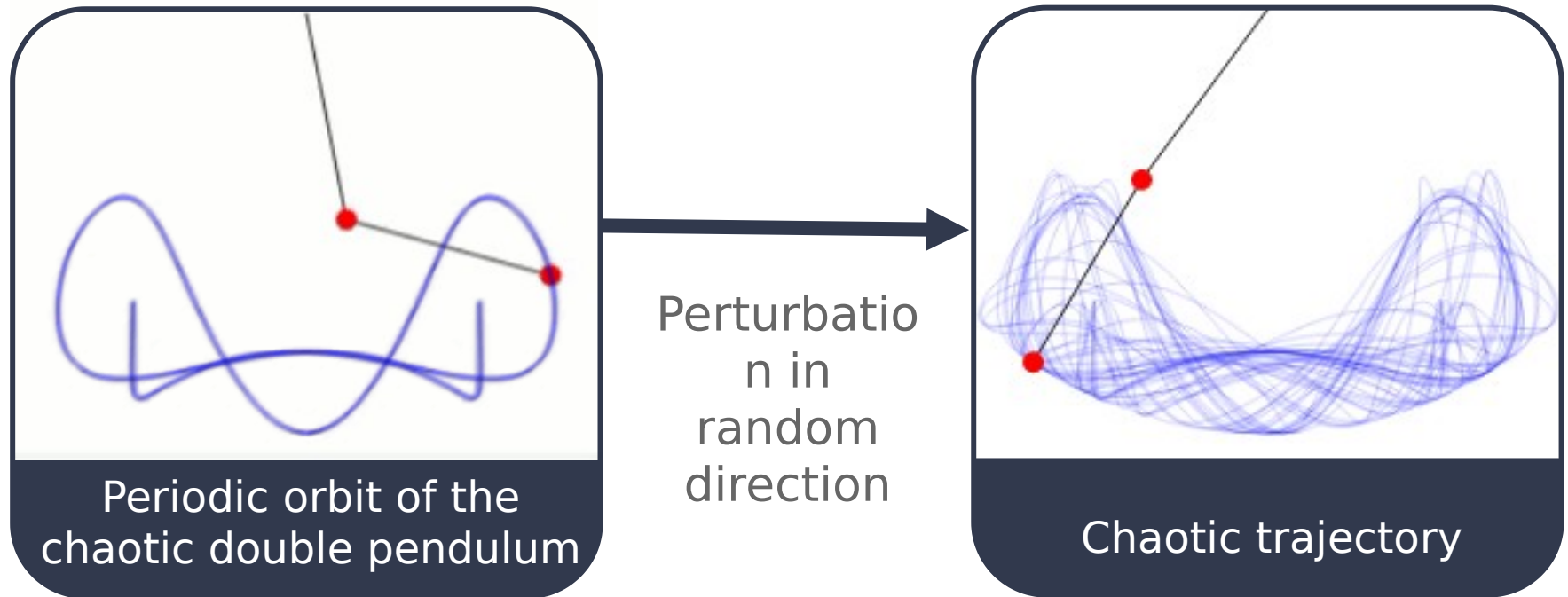
Quantization

Quantum spectrum

Quantum eigenstates



# Chaotic systems also have periodic orbits



# Gutzwiller theory (chaotic systems)

**Unstable** periodic orbits in classical  
limit

Quantization

Quantum spectrum

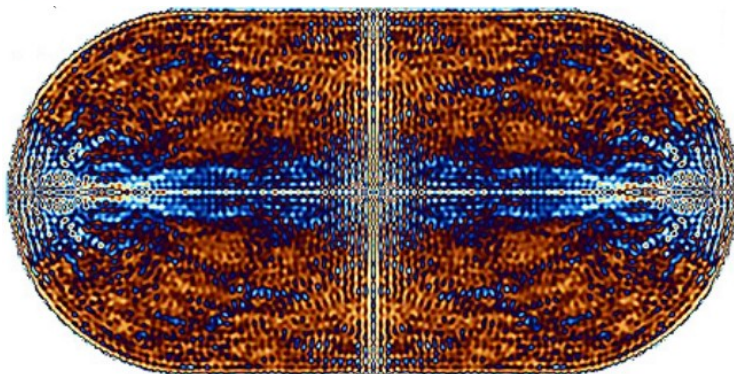
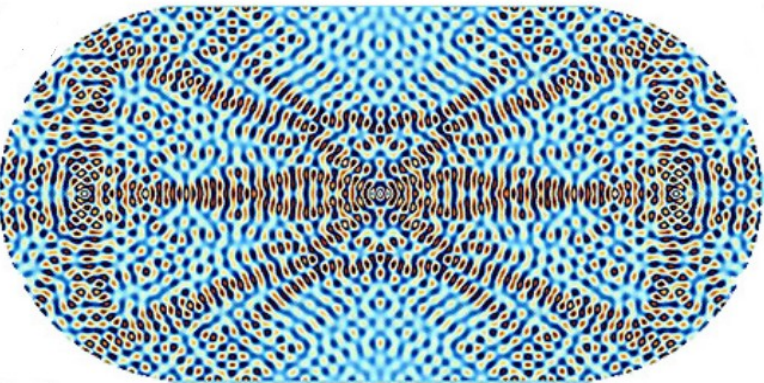
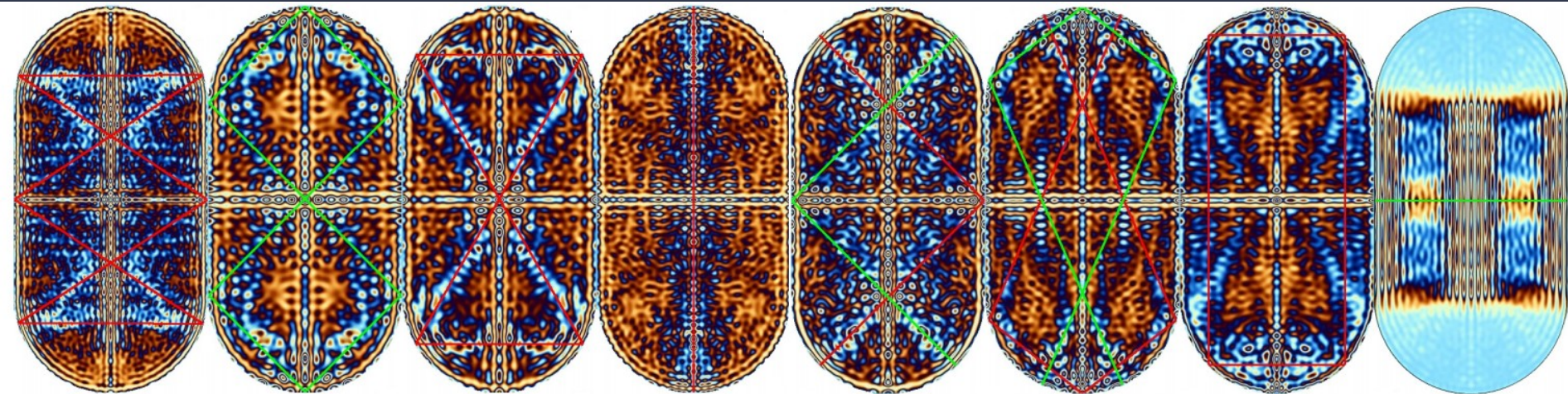
Quantum  
eigenstates?

Yes! (some of  
them)  
-Heller



# Quantum scars

Heller, E. J. Phys. Rev. Lett. **53**,  
1515–1518 (1984)



King, C. C.  
Quanta **3**, 16-  
31 (2014)

# The Dicke model

Dicke, R. H. Phys. Rev. **93**,  
99 (1954).

A **single-mode cavity** of frequency  $\omega$  interacting with **two-level atoms** with level spacing  $\omega_0$ .

$$\hat{H}_D = \underbrace{\omega \hat{a}^\dagger \hat{a}}_{\text{Energy of photons}} + \underbrace{\omega_0 \hat{J}_z}_{\text{Energy of atoms}} + \underbrace{\frac{\gamma}{\sqrt{N}} \hat{J}_x (\hat{a}^\dagger + \hat{a})}_{\text{Coupling}}$$

Energy of photons

Energy of atoms

$$\hat{J}_z = \sum_{i=1}^N \hat{\sigma}_z^{(i)} \quad \hat{J}_y = \sum_{i=1}^N \hat{\sigma}_y^{(i)} \quad \hat{J}_x = \sum_{i=1}^N \hat{\sigma}_x^{(i)}$$

Collective spin operators (all atoms are identical)

Coupling

$\gamma$   
Coupling strength

# The Dicke model: classical limit

- de Aguiar, M. A. M. et al. EPL **15**, 125 (1991),  
- Bastarrachea-Magnani, M. A. et al. PRA **89**, 032101 (2014)

Glauber coherent

$$|q, p\rangle = \hat{D}(q, p)|0\rangle$$

Displacement  
operator

Photon  
vacuum state

Bloch coherent states ( $j = \mathcal{N}/2$ )

$$|Q, P\rangle = \hat{R}(Q, P)|j, -j\rangle$$

Rotation operator

All atoms in  
ground state

Classical Hamiltonian

$$h_{\text{cl}}(\mathbf{x}) = \frac{\langle \mathbf{x} | \hat{H}_D | \mathbf{x} \rangle}{j} =$$

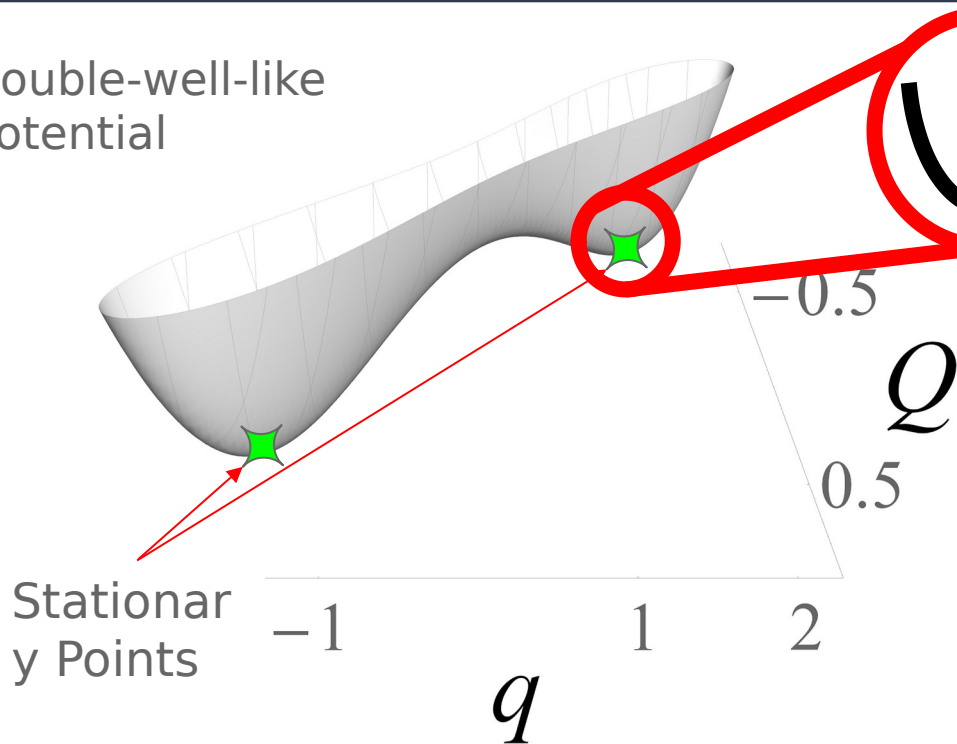
$$\frac{\omega}{2} (q^2 + p^2) + \frac{\omega_0}{2} (Q^2 + P^2) + 2\gamma q Q \sqrt{1 - \frac{Q^2 + P^2}{4}} - \omega_0$$

$$|\mathbf{x}\rangle = |q, p\rangle \otimes |Q, P\rangle$$

$$\mathbf{x} = (q, p, Q, P)$$

# Periodic orbits in the Dicke model

Double-well-like potential



≈ Harmonic oscillator

Normal frequencies

$$\Omega^{A,B}$$

=

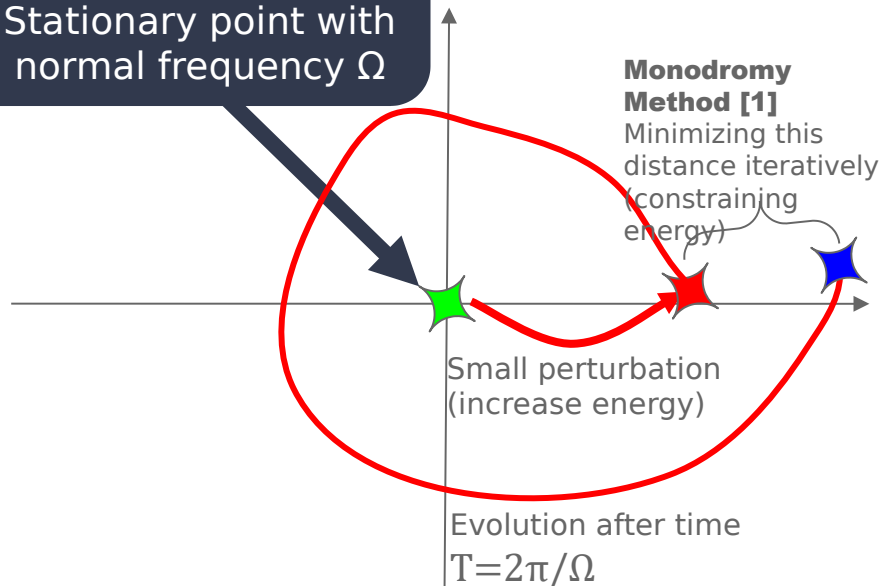
$$\sqrt{\frac{1}{2\omega^2} \left( (16\gamma^4 + \omega^4) \pm \sqrt{(\omega^4 - 16\gamma^4)^2 + 4\omega^6\omega_0^2} \right)}$$

Theorem 2.1 of [Alan Weinstein, Inv. Math. **20**, 47-57 (1973)]

“A family of periodic orbits emanates from  for **each** normal frequency.”

# Finding periodic orbits by perturbing stationary points

Stationary point with normal frequency  $\Omega$



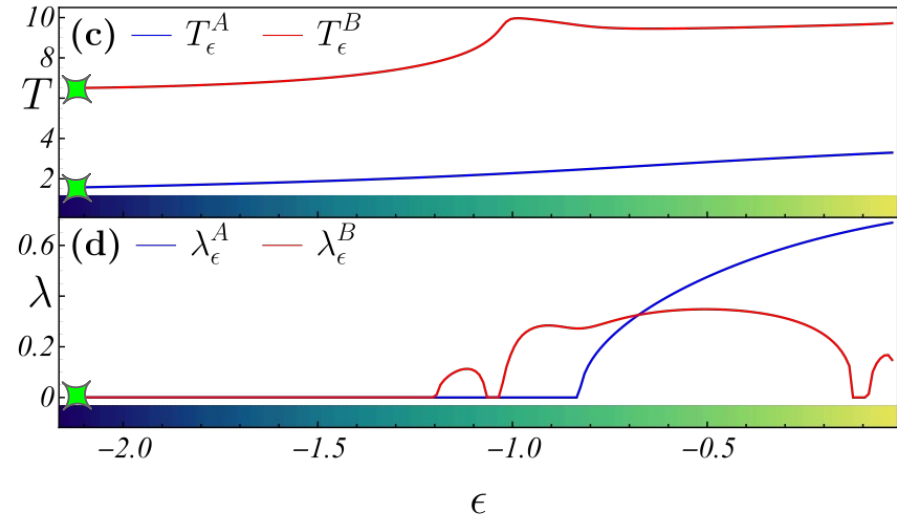
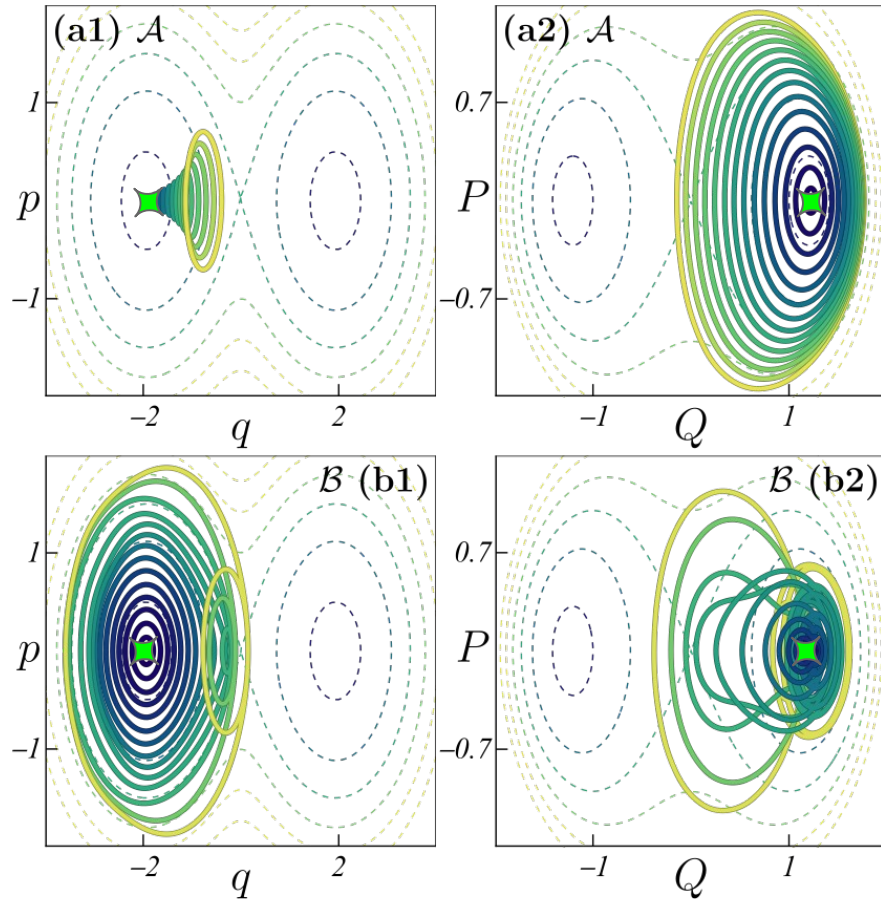
- [1] M.A.M de Aguiar et. al, Ann. Phys. **180** 167-205 (1987)  
[2] S. Pilatowsky Cameo et. al, New. J. Phys. (2021) In press. 10.1088/1367-2630/abd2e6

True periodic orbit with increased energy

repeat

Family of periodic orbits corresponding to fundamental frequency  $\Omega$ . [2]

# The Dicke model: fundamental families of






**New Journal of Physics**

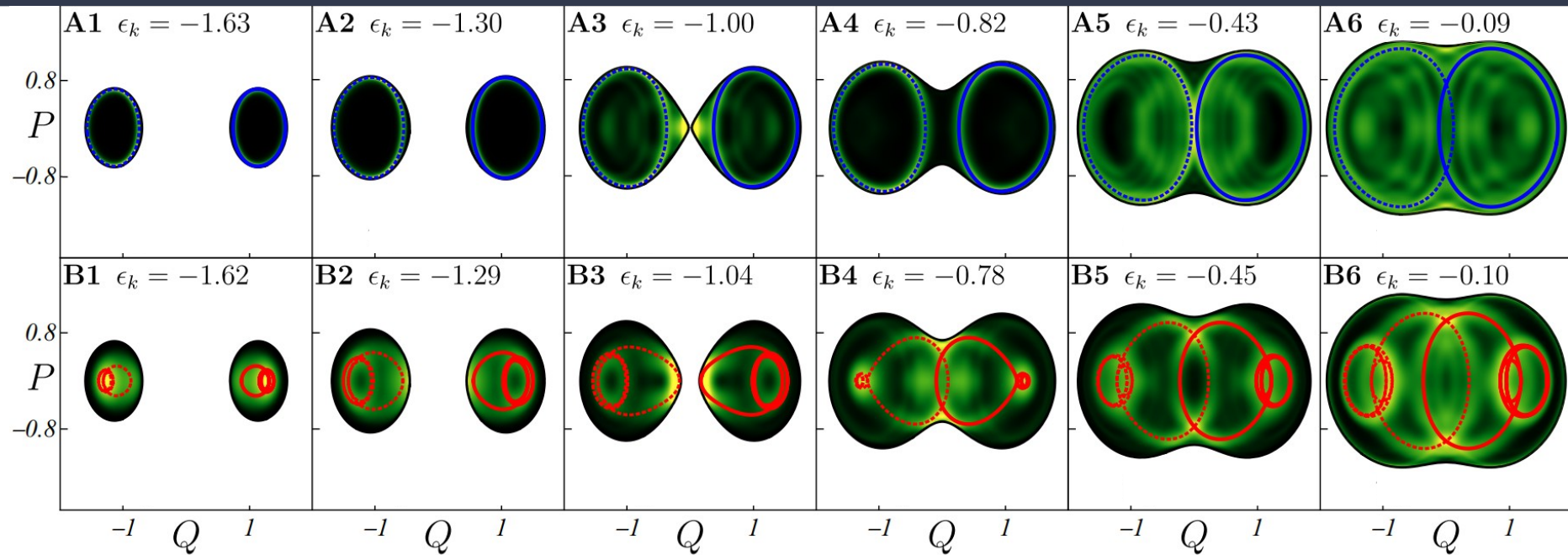
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Quantum scarring in a spin-boson system: fundamental families of periodic orbits

Saúl Pilatowsky-Cameo<sup>1</sup> , David Villaseñor<sup>2</sup>, Miguel A Bastarrachea-Magnani<sup>3</sup>, Sergio Lerma<sup>4</sup>, Lea F Santos<sup>5</sup>  and Jorge G Hirsch<sup>6</sup> 

# The Dicke model: scars of fundamental families



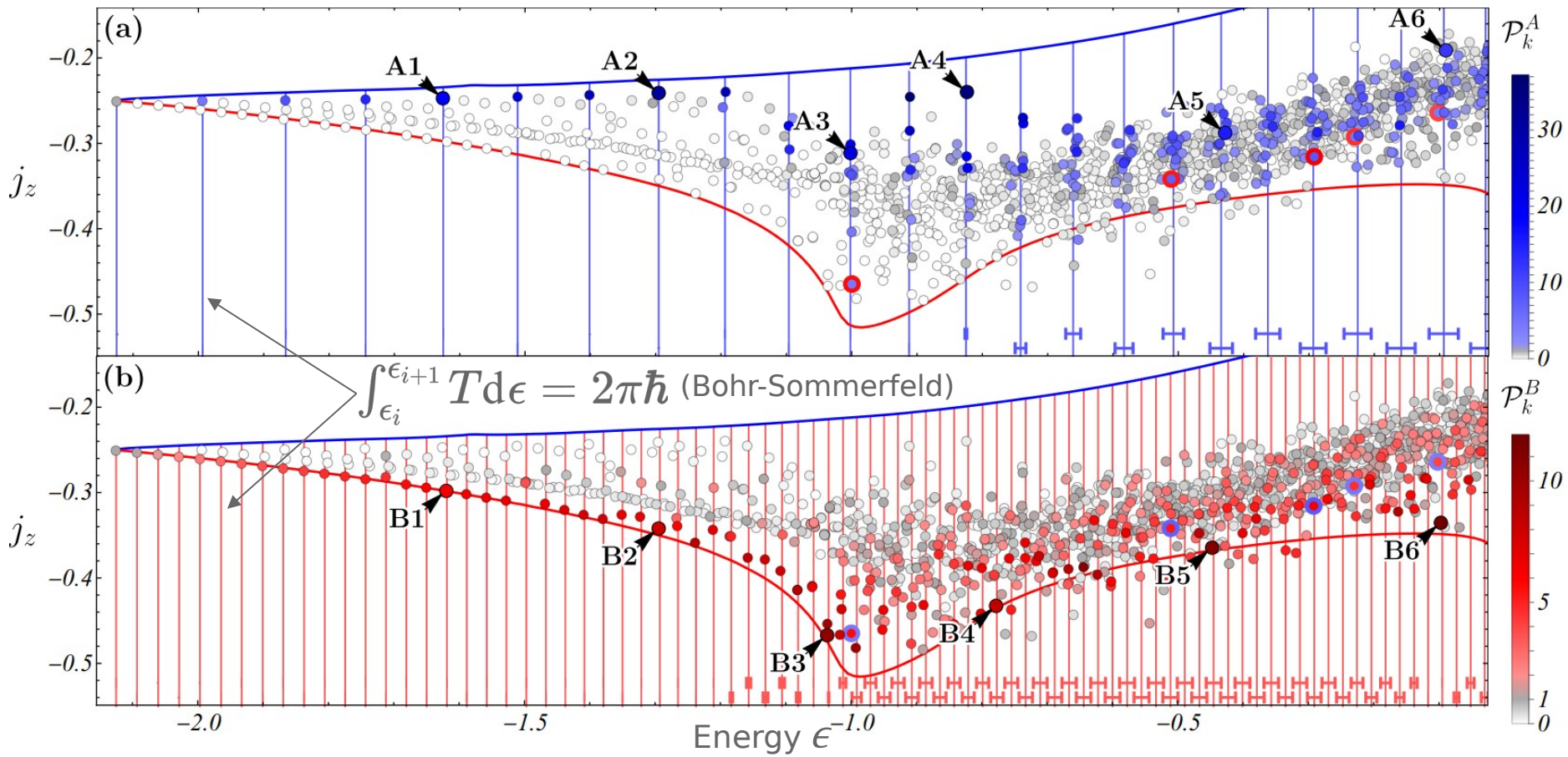
$$\mathcal{P}_k = \frac{\text{tr}(\hat{\rho}_{\text{PO}} \hat{\rho}_k)}{\text{tr}(\hat{\rho}_{\text{PO}} \hat{\rho}_{\text{DS}})}$$

$$\hat{\rho}_k = |E_k\rangle\langle E_k|$$

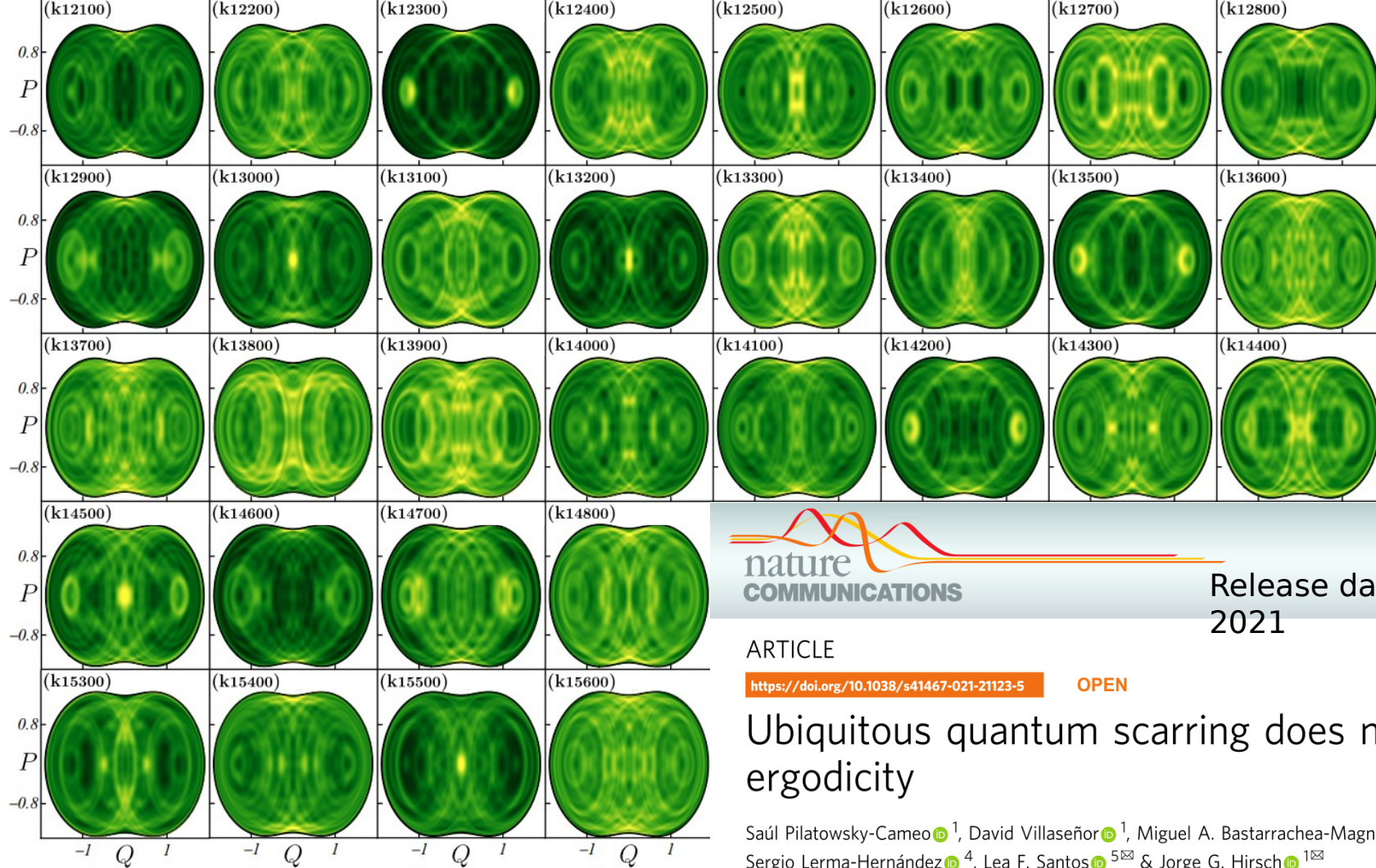
$\hat{\rho}_{\text{PO}}$  = Tubular state around periodic orbit

$\hat{\rho}_{\text{DS}}$  = Totally delocalized state

# Peres lattice colored by scarring measure







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ARTICLE

<https://doi.org/10.1038/s41467-021-21123-5> OPEN

# Ubiquitous quantum scarring does not prevent ergodicity

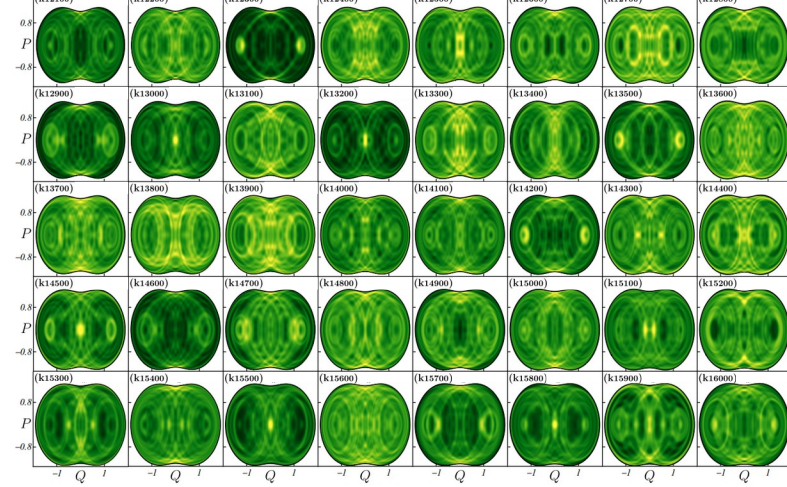
Saúl Pilatowsky-Cameo <sup>1</sup>, David Villaseñor <sup>1</sup>, Miguel A. Bastarrachea-Magnani <sup>2,3</sup>, Sergio Lerma-Hernández <sup>4</sup>, Lea F. Santos <sup>5</sup>✉ & Jorge G. Hirsch <sup>1</sup>✉

## ARTICLE

<https://doi.org/10.1038/s41467-021-21123-5>

OPEN

## Ubiquitous quantum scarring does not prevent ergodicity

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- **All** eigenstates of the chaotic Dicke model are scarred
- Pure states never occupy more than half of the available phase space.
- Quantum ergodicity is achievable only as an ensemble property, after temporal averages are performed

See David Villaseñor's talk tomorrow