

# Overview

## Ergodicity and Chaos

**Sergio Lerma<sup>1</sup> & Lea F. Santos<sup>2</sup>**

*<sup>1</sup>Facultad de Física, Universidad Veracruzana, Veracruz, Mexico*

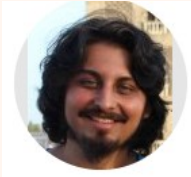
*<sup>2</sup>Department of Physics, Yeshiva University, New York, NY, USA*

Welcoming group

Open to collaborations

Purpose of this overview is not only to promote some final discussions, but to motivate future collaborations with Lerma-Santos-Hirsch's group and among the participants.

# Time line



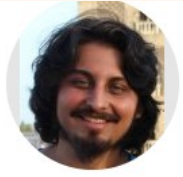
**2011**

## **Researchers and postgraduate students**

Jorge Hirsch

M. Bastarrachea

# Time line



**2011**

**2014**

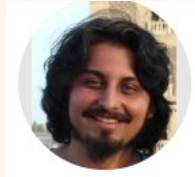
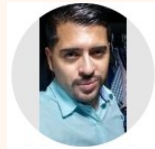
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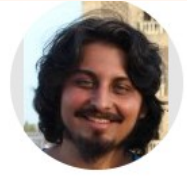
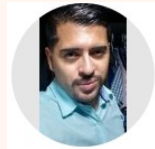
M. Bastarrachea

S. Lerma

J Chávez-Carlos

B. López-del-Carpio

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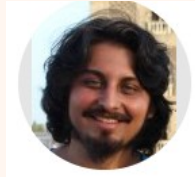
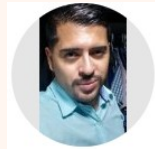
S. Lerma

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B. López-del-Carpio

A. Relaño

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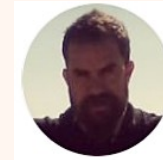
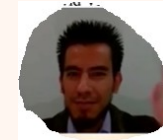
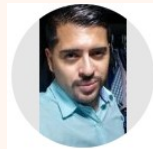
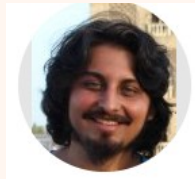
Lea Santos

B. López-del-Carpio

A. Relaño

2018 First Workshop

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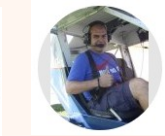
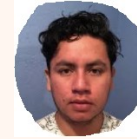
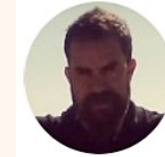
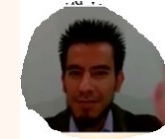
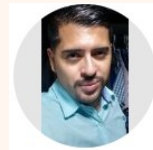
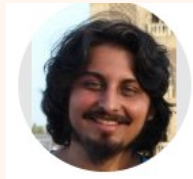
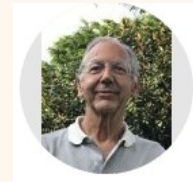


B. López-del-Carpio

A. Relaño

2018 First Workshop  
2019 2nd Workshop

# Time line



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2020

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D. Villaseñor

E.J. Torres

Javier de la Cruz

Saúl-Pilatowsky

B. López-del-Carpio

A. Relaño

Mario Quiroz

Roberto Montiel

2018 First Workshop

2019 2nd Workshop

**2020 3d Workshop**



# Time line



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D.Villaseñor

E.J. Torres

Javier de la

Saúl-Pilatowsky

**Ergodicity and chaos in many-body systems**

Around 50 participants!

A. Reiano

Mario Quiroz

Roberto Montiel

2018 First Workshop

2019 2nd Workshop

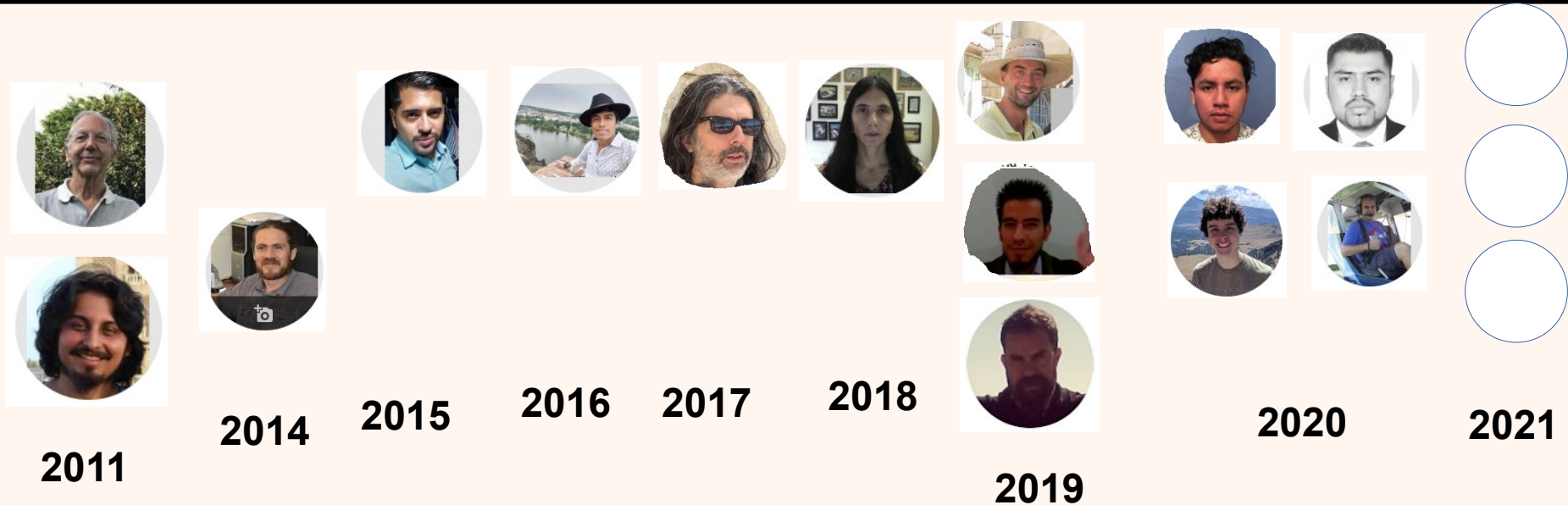
2020 3d Workshop

**2021 4th Workshop**

Ergodicity and chaos, 2021

Overview

# Time line



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E.J. Torres

Javier de la Cruz

Saúl-Pilatowsky

**We hope many new collaborations to come**

B. López-del-Carpio

A. Relaño

Mario Quiroz

Roberto Montiel

2018 First Workshop

2019 2nd Workshop

2020 3d Workshop

**2021 4th Workshop**

# Topics for discussion

## 1) Definitions of ergodicity

- classical ergodicity
  - mixing
- quantum ergodicity
- ETH - chaos
  - self-averaging
  - phase space

## 2) Quantum scarring

Why is it ubiquitous in the Dicke model?

Measures of scarring:

- scar function
- P from NJP
- ETH vs scarring

## 3) Localization

relationship with scarring and lack of ergodicity?

- L (phase space)
  - Robnik's group
  - IPR, R (Hilbert space)
- phase space vs Hilbert space

Eigenvalues vs eigenstates:

- **beta vs PR**
- dependence on basis



ARTICLE

<https://doi.org/10.1038/s41467-021-21123-5>

OPEN

Ubiquitous quantum scarring does not prevent ergodicity

Saúl [Pilatowsky-Cameo](#)<sup>1</sup>, David [Villaseñor](#)<sup>1</sup>, Miguel A. [Bastarrachea-Magnani](#)<sup>2,3</sup>, Sergio [Lerma-Hernández](#)<sup>4</sup>, Lea F. [Santos](#)<sup>5</sup> & Jorge G. [Hirsch](#)<sup>1</sup>

# Topics for discussion

## 5) Wigner-Dyson distribution: What is many-body quantum chaos?

WD needed for thermalization?

WD = chaos in many-body?

WD, but ballistic transport in the SINGLE defect model

## 6) Chaos detectors vs symmetries

- correlation hole
- purity
- off-diagonal elements of observables (SPECK of CHAOS)

Useful also for identifying

- symmetries
- integrable models

## 7) Open systems

- by adding a little bit of **dissipation**,  
find ***stable regions***
- what PR tells us that  
classical phase space does not

# Topics for discussion

## 8) Time scales

quantum-classical correspondence

- Ehrenfest time
- Thouless time (correlation hole)
  
- Diffusion time  $t_D$
- vs
- Heisenberg time  $t_H$
- when  $t_D > t_H$  and  $t_H > t_D$

Relaxation in

real space

momentum space

dependence on observable:

chaos, integrability, initial state

multifractality

## 9) Multifractality

- phase space fragmentation
- vs
- Hilbert space fragmentation

## 10) Correlations in realistic systems

Is there any physical quantity that could detect the correlations of the eigenstates of realistic systems?

Observables:

few-body observable, entropies,  
quantum Fisher, FOTOC, OTOC  
short vs long times

Entropy growth & Kolmogorov-Sinai

# Topics for discussion

## 8) Time scales

quantum-classical correspondence

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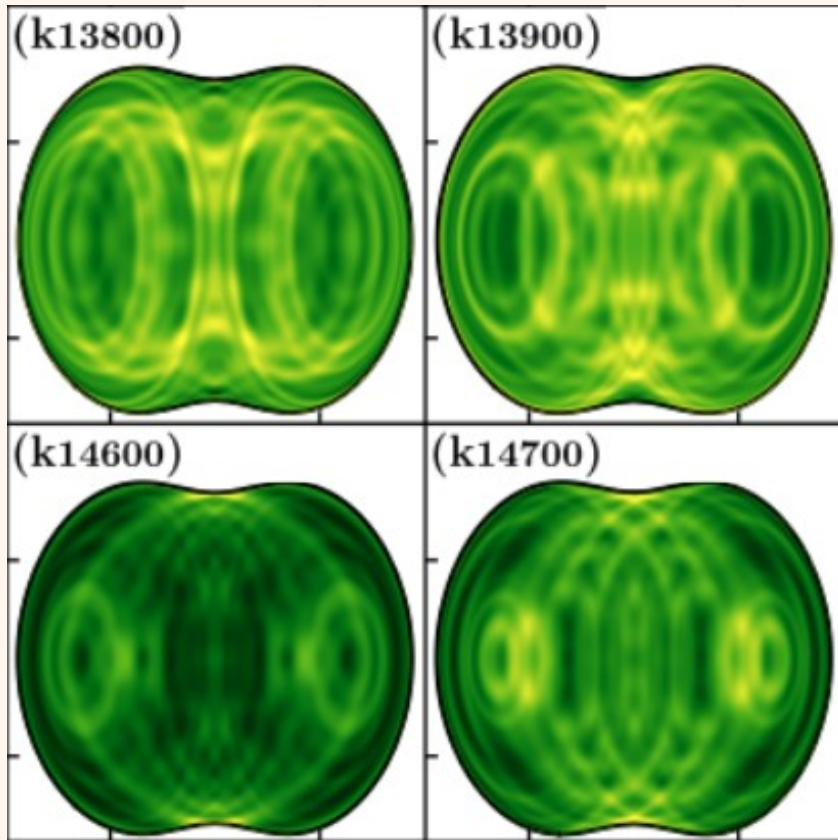
short vs long times

Entropy growth & Kolmogorov-Sinai  
ESQPTs from Wigner and Hus funct

# Topics for discussion

OTHERS????

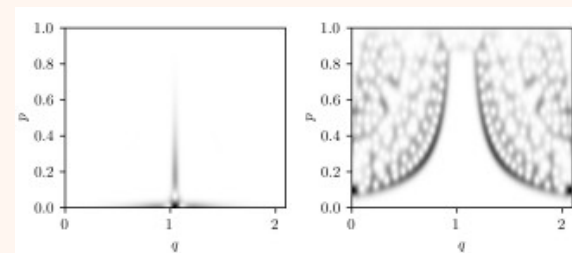
# Quantum scarring



May the semiclassical random wave model presented by Juan Diego help to understand the ubiquitous quantum scarring in the eigenstates of the Dicke model?

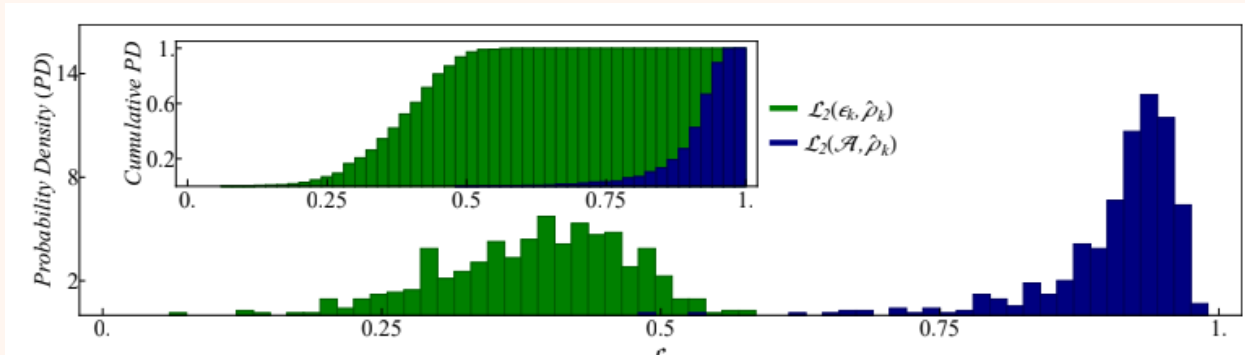
States that are NON-ergodic (localized in phase-space  $L < 1$ ) in phase space lead to a violation of ETH?

Are we missing small regularity island in the large energy region of the Dicke model?





# Localization measures



David Villaseñor  
**L**

Crt Lozej  
Qian Wang  
**A**

How does the tD/tH ratio depends on energy and coupling in the Dicke model?

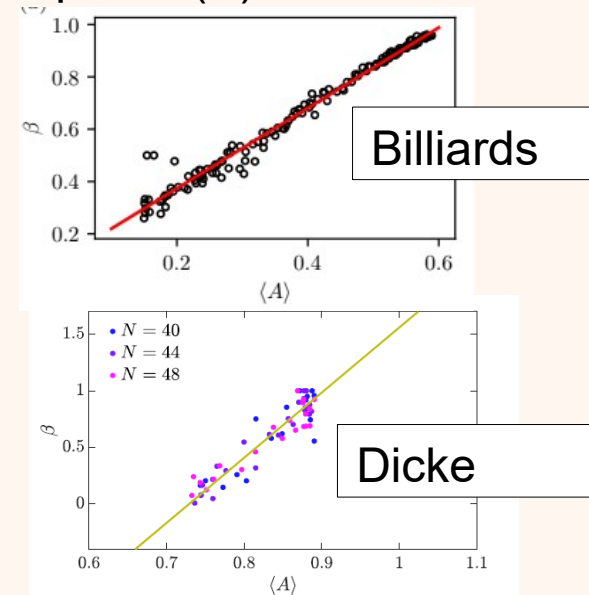
Does the beta distribution describe  $P(L)$  as it does for  $P(A)$ ?

Is the relation between Brody parameter and localization  $L$  also linear?

Localization measures in phase space using

-Projection over energy surfaces (L)

-Projection over Bloch sphere (A)



# Ergodicity

Wikipedia:

The origins of ergodicity lie in [statistical physics](#), where [Ludwig Boltzmann](#) formulated the [ergodic hypothesis](#).

In [physics](#) and [thermodynamics](#), the **ergodic hypothesis** says that, over long periods of time, the time spent by a system in some region of the [phase space](#) of [microstates](#) with the same energy is proportional to the volume of this region, i.e., that all accessible microstates are [equiprobable](#) over a long period of time.

The ergodic hypothesis is often assumed in the [statistical analysis](#) of [computational physics](#). The analyst would assume that the **average of a process parameter over time and the average over the statistical ensemble are the same.**

This assumption - that it is as good to simulate a system over a long time as it is to make many independent realizations of the same system – is not always correct. (FPUT)

# Quantum Ergodicity vs ETH

$$\langle O(t) \rangle = \langle \Psi(t) | O | \Psi(t) \rangle = \sum_{\alpha \neq \beta} C_{\beta}^{ini*} C_{\alpha}^{ini} e^{i(E_{\beta} - E_{\alpha})t} O_{\beta\alpha} + \sum_{\alpha} |C_{\alpha}^{ini}|^2 O_{\alpha\alpha}$$

$$|\Psi(0)\rangle = \sum_{\alpha} C_{\alpha}^{ini} |\alpha\rangle$$

$$O_{\beta\alpha} = \langle \beta | O | \alpha \rangle$$

$$O_{\alpha\alpha} = \langle \alpha | O | \alpha \rangle$$

Infinite time average

Thermodynamic average

$$\overline{\langle O(t) \rangle} \equiv \sum_{\alpha} |C_{\alpha}^{ini}|^2 O_{\alpha\alpha} \xrightarrow{=?} O_{micro} \equiv \frac{1}{N_{E_0, \Delta E}} \sum_{\alpha} O_{\alpha\alpha}$$

depends on the initial conditions

depends only on the energy

ETH: the expectation values  $O_{\alpha\alpha}$  of few-body observables do not fluctuate for eigenstates close in energy

**Chaos**  
guarantees  
thermalization, ETH

- **CHAOTIC STATES**
- **FEW-BODY O**
- **SELF-AVERAGING**

**Chaotic states**

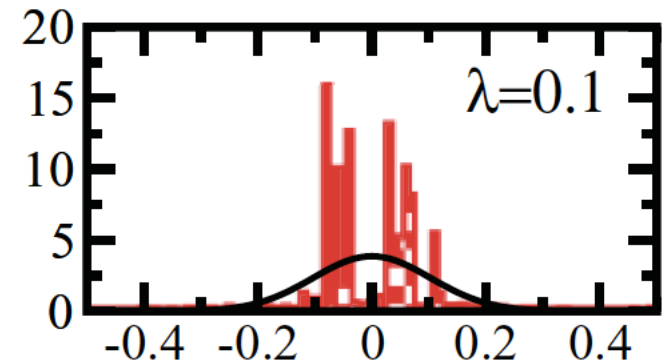
# Chaotic Eigenstates

$$H = H_0 + \lambda V$$

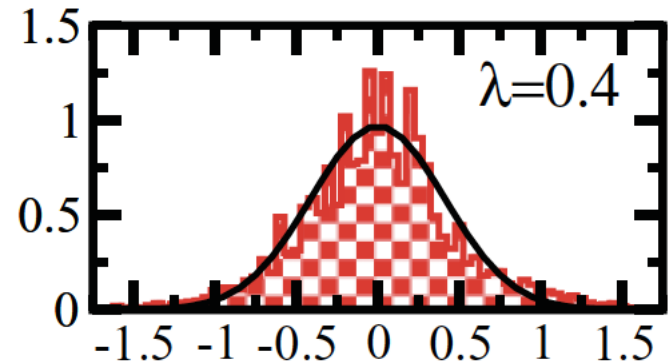
- **CHAOTIC STATES**  
*Up to which point?*
- **FEW-BODY Observables**

Chaos  
guarantees  
thermalization, ETH

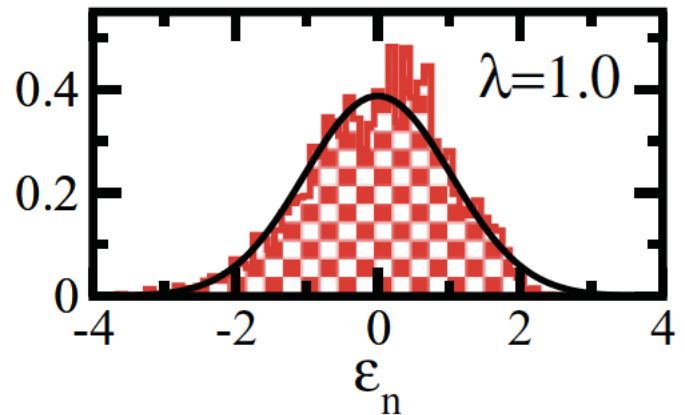
$|C_n^{\alpha}|^2$



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# Quantum Ergodicity vs ETH

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Infinite time average

Thermodynamic average

$$\overline{\langle O(t) \rangle} \equiv \sum_{\alpha} |C_{\alpha}^{ini}|^2 O_{\alpha\alpha} \xleftrightarrow{=} O_{micro} \equiv \frac{1}{N_{E_0, \Delta E}} \sum_{\substack{\alpha \\ |E_0 - E_{\alpha}| < \Delta E}} O_{\alpha\alpha}$$

depends on the initial conditions

depends only on the energy

Literature of many-body systems:

~~Chaotic regime, chaotic states~~

Thermal phase or ergodic phase, ergodic states

Why this preference?

What is

many-body quantum chaos?  
(lack of semiclassical)

--WD? -- Diffusive transport?

**Eigenstates**

# Quantum Ergodicity vs ETH

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QUESTIONS:

- Dependence on initial state?
- Ergodic state! Too strong?
- Hilbert space
- Dependence on observables

# Quantum Ergodicity vs Phase Space

In the classical limit, a system is ergodic if the trajectories cover the energy shell homogeneously. We then adopt the same definition for quantum ergodicity. To quantify how much of the energy shell is visited on average by the evolved state

$$\hat{\rho}(t) = e^{-i\hat{H}_D t} \hat{\rho} e^{i\hat{H}_D t}$$

we consider the infinite-time average

$$\bar{\rho} = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T dt \hat{\rho}(t)$$

and compute

$$\bar{\mathcal{L}}(\epsilon, \hat{\rho}) \equiv \mathcal{L}(\epsilon, \bar{\rho})$$

Measure of localization in phase space

$$\mathcal{L}(\epsilon, \hat{\rho})^{-1} = \frac{1}{N} \int_{\mathcal{M}_\epsilon} ds Q_{\hat{\rho}}^2(\mathbf{x})$$

If the whole energy shell is homogeneously visited by  $\hat{\rho}$ , then

and the  $\bar{\mathcal{L}}(\epsilon, \hat{\rho}) = 1/\epsilon$  is ergodic.

**Stationary states (eigenstates of Dicke, superpositions) are NON-ergodic**

**Non-stationary states can be: Random states and coherent states**

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- ❖ Quantum ergodicity in phase space vs quantum ergodicity in Hilbert space (ETH)?
- ❖ Eigenstates that are NON-ergodic in phase space lead to a violation of ETH?  
*no -- few-body observables*

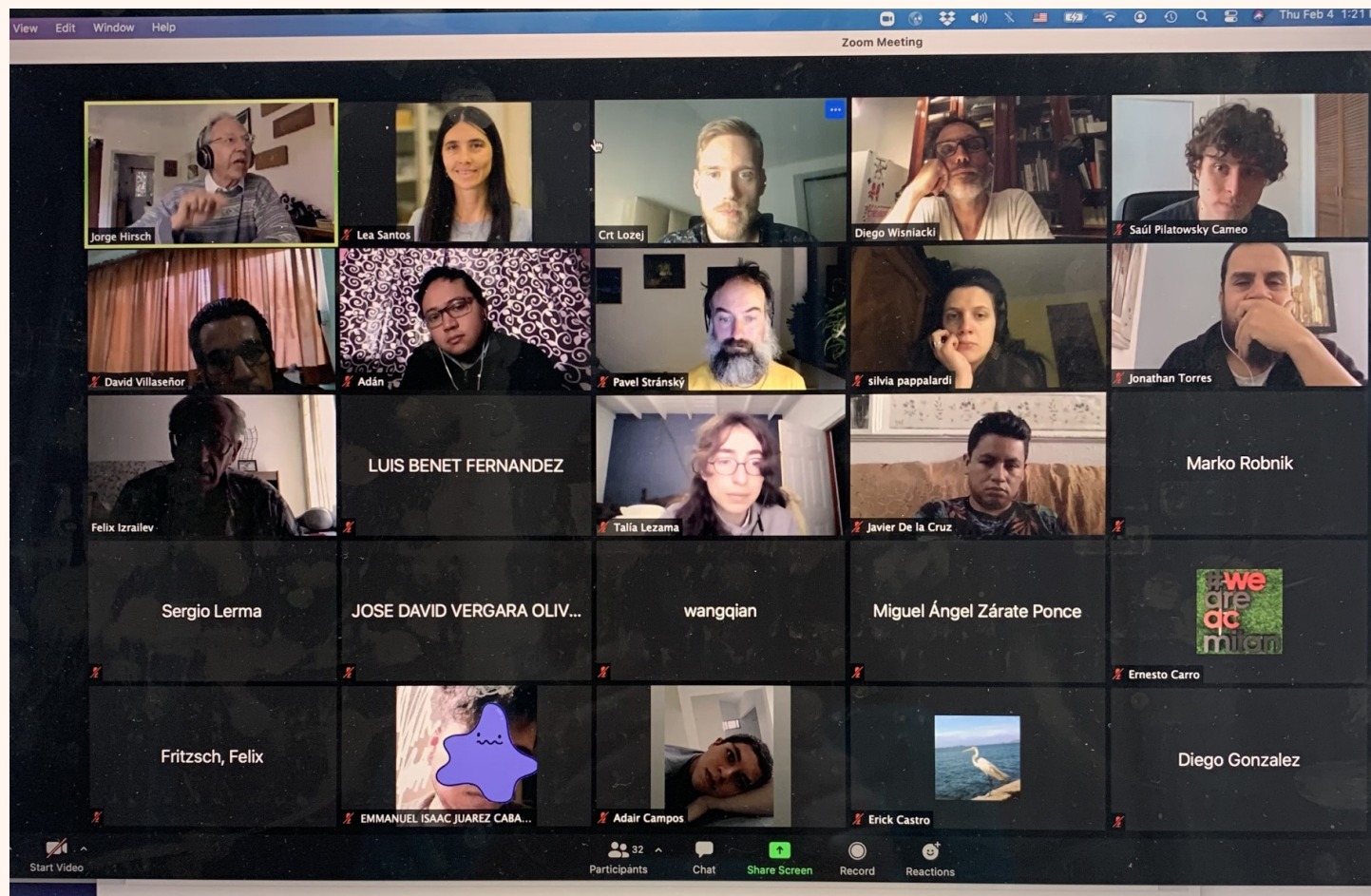




# Ergodicity and chaos in many-body systems

Around 50 participants!

2021 4th Workshop



Ergodicity and chaos, 2021

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