Introduction to Higgs Physics (3rd Lesson)

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September 27, 2016

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- **2** THE Standard Model
- **3** SM Higgs Properties and its detection at LHC
- 4 Higgs Hysics Beyond the SM
- **5** Higgs Physics in the far UV

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Then what?



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What is the nature of EWSB?

Questions:

- Is there only one Higgs doublet that generates the masses of all particles?
- **2** Will we be able to test Higgs couplings with light fermions?
- **3** Are the Higgs couplings diagonal in flavor space?
- Why W-mass << Planck mass? ((Hierarchy problem))

Possible Answers:

- Strongly Interacting -Higgsless world DECEASED!
- Strongly Interacting Composite Higgs pNGB,
- Weakly interacting- SM valid up to Planck Scale,
- Weakly interacting- Multi-Higgs model (SUSY, THDM, etc),

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Higgs identity: $g_{hXX} = c_X g_{hXX}^{sm}$

In the SM: $c_X = 1$,

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The Universal Higgs fit - P. Giardino et al., arXiv:1303.3570 [hep-ph]

Under the small deviations approximation:

$$c_X = (1 + \epsilon_X) \tag{1}$$

From a fit to all observables (signal strengths), and assuming no new particles contribute to the loop decays hgg and $h\gamma\gamma$, they get:

- hZZ (hWW): $\epsilon_Z = -0.01 \pm 0.13$ ($\epsilon_W = -0.15 \pm 0.14$),
- *hbb*: $\epsilon_b = -0.19 \pm 0.3$,
- $h\tau\tau: \epsilon_{\tau} = 0 \pm 0.18$
- *htt* (from *hgg*): $\epsilon_t = -0.21 \pm 0.23$

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SM Higgs identity: $g_{hXX}^{sm} = \frac{M_X}{v}$



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Higgs Couplings in 3+1 HDM



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Higgs parity: $h\bar{f}f$ or $h\bar{f}\gamma_5 f$?



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Higgs parity: $h\bar{f}f$ or $h\bar{f}\gamma_5 f$?



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Higgs and Flavor Violation Only FC Higgs couplings $(h\bar{f}f)$ or also possible FV $(h\bar{f}_if_j)$?, ex. $h \to \tau \mu$ (Diaz-Cruz and collab.),

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2014/07/05

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Search for lepton flavor violating decays of the Higgs boson



The CMS Collaboration

The Hierachy problem

When an scalar interacts with a heavy fermion M, with $L_Y = y \bar{\Psi} \Psi \phi$, and UV cutoff Λ , the scalar mass gest corrected, i.e.

$$m_h^2 = m_0^2 + \frac{y^2}{16\pi^2} [c_1 \Lambda^2 + c_2 m_0^2 ln \frac{\Lambda}{m} + M^2]$$
(2)

The problem: $m_h = 125 - 126 \text{ GeV}$ but since $\Lambda >> O(1)$ TeV, need a large cancellation.

Some solutions:

- Composite Higgs (as in QCD!),
- Higgs is part of D dim vector field: $A_M = (A_\mu, A_i)$,
- Cancelation between boson-fermion loops (\rightarrow SUSY),
- Accidental cancelacion (Veltman's condition):

$$\lambda = y_t^2 - \frac{1}{8} [3g^2 + g'^2] \tag{3}$$

NO LONGER WORKS!...at the EW scale ($\rightarrow m_h \simeq 200 \text{ GeV}$,)

....BUT WHAT ABOUT AT M_{pl} ?

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Other problems in the SM

- Large/Little hierarchy problem,
- Neutrino masses and flavor problem,
- Strong CP problem,
- Dark Matter,
- Cosmological constant (Dark energy),
- Some deviations from the SM (a few std. dev.), e.g. Δa_{μ} , etc.
- Aesthetical questions,

They all suggest the need for New Physics.

Beyond the SM

- Models with new fermions (4ta family, etc)
- Models with new gauge forces (U(1)', Left-Right, ..)
- Models with extra Higgs multiplets (2HDM, triplets,..)
- Models with Grand Unification (ex. $SU(5), SO(10), E_6,...$)
- Models with new symmetries (SUSY),
- Models with extra dimensions extra.
- etc.

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A modern view of Physics BSM

Physics BSM incorporates Extra Dimensions,

 Fermionic XDx^µ → Z^M = (x^µ, Q, Q̄) : Supersymmetry
 Bosonic XD-

$$x^{\mu} \to X^M = (X^{\mu}, X^i)$$
: Large Extra Dimensions,

• Curved XD $x^{\mu} \to X^{M} = (X^{\mu}, X^{i})$: Randall-Sundrum

AdS/CFT duality means XD \rightarrow Strong Ints.



Supersymmetry (SUSY)

Why is SUSY attractive? It is a new symmetry that relates fermions and bosons,

- Offers the possibility to stabilize the Higgs mass and EWSB,
- Improves Unification and o.k. with proton decay,
- Favors a light Higgs boson, in agreement with EWPT (and LHC?), i.e. $m_h \leq 160$ GeV,
- New sources of flavor and CP violation may help to get the right BAU,
- LSP is stable and a possible Dark matter candidate.

Gauge Coupling Unification



The MSSM

The minimal extension of the SM consistent with SUSY, is based on:

- SM Gauge Group (\rightarrow gauge bosons and gauginos),
- 3 families of fermions and sfermions,
- Two Higgs doublets $(H_u \text{ and } H_d)$,
- Soft-breaking of SUSY (Hidden sector),
- R-parity distinguish SM and their superpartners \rightarrow LSP is stable and DM candidate.

The MSSM particle content

	SM	Superpartners		
SM	W^{\pm}, Z, γ	Wino,Zino, Photino		
Bosons	gluon	gluino		
	Higgs bosons	Higgsinos		
SM	quarks	squarks		
Fermions	leptons	sleptons		
	neutrinos	$\operatorname{sneutrinos}$		

Mixing of gauginos and Higgsinos \rightarrow Charginos (χ_i^{\pm} , i = 1, 2) and Neutralinos (χ_j^0 , j = 1, 4),

Gravitino is also part of the spectrum.

The parameters of the MSSM

In addition to SM parameters, the MSSM includes $\mathrm{O}(100)$ new ones:

- Scalar masses (Sleptons, squarks, Higgs),
- Gaugino masses $(\tilde{M}_G, \tilde{M}_W, \tilde{M}_B)$,
- Trilinear terms $(A_{\tilde{f}} \text{ for squarks and sleptons}),$
- From Higgs sector: $\tan \beta = v_2/v_1$ and μ ,
- The masses of superpartners have important implications for EWSB,
- Spectrum of superpartners depends on mechanism of SUSY breaking,

Susy Spectrum



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MSSM Higgs Potential

At tree-level MSSM Higgs sector is a 2HDM of type-II, i.e. it contains two Higgs doublets, with Potential:

Lagrangian. The F terms contribute

$$V_F = \mu^2 (H_u^{0*} H_u^0 + H_d^{0*} H_d^0)$$

The D terms contribute

$$V_D = \frac{g^2 + g'^2}{8} (H_u^{0*} H_u^0 - H_d^{0*} H_d^0)^2$$

The soft SUSY breaking terms contribute

$$V_{soft} = M_{Hu}^2 H_u^{0*} H_u^0 + M_{Hd}^2 H_d^{0*} H_d^0 - (B\mu H_u^0 H_d^0 + h.c.)$$

The MSSM Higgs spectrum

- CP-even neutral Higgs bosons h^0, H^0 , at tree-level $m_h < m_Z$,
- CP-odd neutral Higgs A^0 with $m_H^2 = m_A^2 + m_Z^2 \sin^2 2\beta$,
- Charged Higgs H^{\pm} , with $m_{H^+}^2 = m_A^2 + m_W^2$,
- Masses and mixing angles fixed with: m_A and $tan\beta = v_2/v_1$,
- When $m_A \leq \tilde{m}$, Higgs search uses SM techniques.
- But H^0, A^0, H^{\pm} may decay into SUSY modes; LHC search gets more complicated!,

The MSSM Higgs mass

Radiative effects of Stop-top loops can make: $m_h > m_Z$

$$m_h^2 = m_Z^2 \left[1 + \frac{3m_t^2}{2\pi^2 m_Z^2} log(\frac{m_{stop}}{m_t})\right]$$
(4)

But to get $m_h = 125$ GeV, with SM-like couplings, need:

- Large superpartner masses O(1) TeV,
- Only a few superpartners could be at the reach of LHC,
- Split SUSY? High Scale SUSY?
- O(1) or large $tan\beta$ allowed,
- Large $tan\beta \rightarrow$ enhanced production of H + bb at LHC,

MSSM Higgs mass (Giudice and Strumia)

SPLIT SUSY



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MSSM Higgs couplings:

$$\begin{array}{ll} \bullet \ (hVV): & \frac{2m_V^2}{v}\cos(\beta-\alpha), \quad v^2=v_1^2+v_2^2, \\ \bullet \ (huu): & \frac{m_u}{v}(\frac{\cos\alpha}{\sin\beta}), \\ \bullet \ (hdd): & \frac{m_d}{v}(\frac{\sin\alpha}{\cos\beta}), \\ \bullet \ (hll): & \frac{m_l}{v}(\frac{\sin\alpha}{\cos\beta}), \\ \bullet \ (hhh): & \simeq \lambda v, \quad \lambda=\frac{g^2+{g'}^2}{8}, \\ \bullet \ (hhhh): & \simeq \lambda. \end{array}$$

Similar expressions hold for H^0, A^0 and H^{\pm} .

Heavy Higgses at LHC



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Composite PGB Higgs



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Composite Higgs (A. Pomarol, ICHEP12)

The light Higgs can be a kind of pion from a new strong sector

The spectrum of the new strong sector could be:



Higgs mass and new physics (A. Pomarol, ICHEP12)

Higgs mass range



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LHC tests of the SM:



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Implications for Dark Matter



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Observed

MSSM Higgs and Dark matter

For heavy sfermions the DM relic density is:

$$\Omega_X h^2 = C_X \left(\frac{m_X}{TeV}\right)^2 \tag{5}$$

- For DM X = pure Bino, no aceptable solution,
- For DM $X = \tilde{H}$ pure Higgsino, $C_{\tilde{H}} = 0.09$ and an aceptable solution is obtained for $1 < M_{\tilde{H}} < 1.2$ TeV,
- For DM $X = \tilde{W}$ pure Wino, $C_{\tilde{H}} = 0.02$ and an aceptable solution is otained for $2 < M_{\tilde{W}} < 2.5$ TeV,

In such case detection at LHC may be harder,

LSP Composition



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Holographic Dark matter

- Composite Higgs can have a "baryon" partner,
- This composite state can be (Holographic) Dark matter (J.L. Diaz-Cruz, PRL81, 2008),
- Deviations from SM Higgs properties can show evidence of dark matter,

The Higgs and the roots of Physics Was it premature to rule our Veltman condition?

A special Value of λ at M_{planck} ?

ML '86



downward flow of RG trajectories \Rightarrow IR QFP \Rightarrow random λ flows to $m_{\rm H} > 150 \ GeV$ $\Rightarrow m_{\rm H} \simeq 126 \ GeV$ flows to tiny values at $M_{\rm Planck}...$

Holthausen, ML Lim (2011) Different conceivable special conditions:

- Vacuum stability $\lambda(M_{pl}) = 0$ [7–12]
- vanishing of the beta function of λ $\beta_{\lambda}(M_{pl}) = 0$ [9, 10]
- \bullet the Veltman condition [13–15] ${\rm Str} {\cal M}^2~=~0,$

$$\begin{split} \delta m^2 &= \frac{\Lambda^2}{32\pi^2 v^2} \mathrm{Str} \mathcal{M}^2 \\ &= \frac{1}{32\pi^2} \left(\frac{9}{4} g_2^2 + \frac{3}{4} g_1^2 + 6\lambda - 6\lambda_t^2 \right) \Lambda^2 \end{split}$$

• vanishing anomalous dimension of the Higgs mass parameter $\gamma_m(M_{pl})=0,\ m(M_{pl})\neq 0$

M. Lindner, MPIK (BUAP) SCALARS 2013, Warsaw Introduction to Higgs Physics (3

Hows does the Higgs potential looks at higher energies?



From JR Espinosa et al.



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Higgs mood



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Why I believed in the Higgs and BSM

Is the Higgs something natural? I would say, yes.

Spin and Isospin:

T / S	0	1/2	1	3/2	2
0	?	Neutrinos-R	gluon	?	?
1/2	Higgs	electron	?	?	?
		quarks			
1	?	?	W, Z	?	?

$$Q_{em} = T_3 + Y \tag{6}$$

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(Where have all the large representations gone?)

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Conclusions.

- LHC is already giving great results,
- Evidence for a SM-like Higgs with $m_h = 125 \text{ GeV}$,
- No evidence at LHC, so far, of new phyics,
- Still possible to find evidence of Dark matter,
- Tests of Higgs couplings at LHC could show deviations from SM (3+1 HDM),
- FCNC decays of Higgs/top could also provide another window into PBSM,
- If no signal of BSM physics shows up at LHC, then what? Super-split SUSY

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Fronteras del micro y macro cosmos



Ademas, podemos agregar una "Frontera Conceptual" (JLDC),



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"This could be heave or this could be hell .."



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Interesting times!



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